

The Fokker-Planck Equation

Session 22 · How distributions of LP states evolve through time

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What we'll cover today

1

Forward Kolmogorov / Fokker-Planck

Distribution dynamics

2

Derivation intuition

Conservation of probability

3

Stationary solutions

Long-run distribution

4

GE-LAV application

Cross-sectional LP state distribution

5

Numerical methods preview

Finite-difference solvers

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Fokker-Planck: how probability flows

For SDE $dL = \mu(L,t)dt + \sigma(L,t)dW$, the density $p(L,t)$ satisfies:

$$\frac{\partial p}{\partial t} = -\frac{\partial}{\partial L} [\mu(L, t) p] + \frac{1}{2} \frac{\partial^2}{\partial L^2} [\sigma^2(L, t) p]$$

Two terms with physical meaning:

Advection: $-\partial/\partial L [\mu(L,t) \cdot p]$

How the drift pulls probability mass along the L axis.

Diffusion: $\frac{1}{2} \cdot \partial^2/\partial L^2 [\sigma^2(L,t) \cdot p]$

How random shocks spread out the probability mass.

What the FP equation describes

The forward partial differential equation for a density.

Given an SDE

$$dX = \mu(X,t) \cdot dt + \sigma(X,t) \cdot dW$$

Question

How does the density $p(X,t)$ of X_t evolve over time?

Answer

Fokker-Planck equation (also called forward Kolmogorov)

Discrete analog

Markov chain master equation

Why 'forward'

Evolves p forward in time from initial $p(x,0)$

Companion 'backward'

Backward equation for value functions (HJB)

Session 22 summary

What we accomplished today

1 Fokker-Planck (forward Kolmogorov) describes how SDE densities evolve through time

2 Two terms: advection (drift) + diffusion (noise spreading)

3 OU process has explicit stationary solution: $\text{Normal}(\bar{L}, \sigma^2/2\kappa)$

4 GE-LAV uses Fokker-Planck for the cross-sectional LP distribution evolution

Next session

Session 23: Jensen's inequality and convexity — the math behind Failure 2

The FP equation: form

Standard form for an Itô diffusion.

PDE

$$\partial p / \partial t = -\partial / \partial X [\mu(X, t) \cdot p] + \frac{1}{2} \cdot \partial^2 / \partial X^2 [\sigma^2(X, t) \cdot p]$$

First term

Drift convection $\cdot -\partial / \partial X [\mu \cdot p]$

Second term

Diffusion $\cdot \frac{1}{2} \cdot \partial^2 / \partial X^2 [\sigma^2 \cdot p]$

Conservation

$\int p \, dX = 1$ preserved (probability is conserved)

Initial condition

$p(X, 0) = p_0(X)$ \cdot prior distribution

Boundary conditions

$p \rightarrow 0$ at infinity; absorbing or reflecting at finite boundaries

Derivation: via test functions

How to derive FP rigorously.

Step 1

Take an arbitrary smooth test function $f(X)$

Step 2

Compute $d/dt E[f(X_t)]$ using Itô

Step 3

Express as $\int f \cdot \partial p / \partial t dX$ (one side) = $\int (\mu \cdot f' + \frac{1}{2} \sigma^2 \cdot f'')$ · $p dX$ (other)

Step 4

Integrate by parts (twice on diffusion term)

Step 5

Use vanishing boundary terms (decay at infinity)

Step 6

Since f arbitrary, equate integrands → FP equation

FP for the OU process

Specialize to our case.

OU SDE

$$dL = \kappa(\bar{L} - L) \cdot dt + \sigma \cdot dW$$

FP equation

$$\partial p / \partial t = \kappa \cdot \partial / \partial L [(L - \bar{L}) \cdot p] + \frac{1}{2} \sigma^2 \cdot \partial^2 p / \partial L^2$$

Stationary solution

$$\partial p / \partial t = 0 \rightarrow p_{\infty}(L) = \sqrt{(\kappa / \pi \sigma^2)} \cdot \exp(-\kappa(L - \bar{L})^2 / \sigma^2)$$

Recognize

Normal(\bar{L} , $\sigma^2 / (2\kappa)$) — matches our prior derivation

Time-dependent

$p(L, t) \rightarrow$ Normal with mean $(L_0 - \bar{L}) \cdot e^{-\kappa t} + \bar{L}$

Variance grows

$$(\sigma^2 / (2\kappa)) \cdot (1 - e^{-2\kappa t})$$

Stationary distribution: GE-LAV calibration

Numerical values for our parameters.

Parameter	Value	Interpretation
κ	0.45/yr	Mean reversion speed
σ	0.32	Diffusion volatility
\bar{L}	1.0	Long-run mean state
Stationary std	0.337	$= \sqrt{(\sigma^2/2\kappa)} \approx \sqrt{(0.32^2/0.9)}$
90% CI for L	[0.45, 1.55]	Most of distribution in normal range
$\Pr(L < -0.5)$	0.07%	Tail probability
$\Pr(L < -1.5)$	<0.001%	Effectively zero under stationary OU

Pr(crisis) under stationary OU is tiny — explains why we see crises rarely.

Spectral decomposition of FP operator

Eigenvalue analysis · convergence rates.

FP operator L_{FP}

$$L_{FP} \cdot p = -\partial/\partial L[\mu \cdot p] + \frac{1}{2} \cdot \partial^2/\partial L^2[\sigma^2 \cdot p]$$

Eigenvalue equation

$$L_{FP} \cdot \varphi_n = -\lambda_n \cdot \varphi_n$$

Smallest eigenvalue

$$\lambda_0 = 0 \cdot \text{stationary state}$$

Second eigenvalue

$$\lambda_1 = \kappa \cdot \text{the spectral gap}$$

Implication

$$\text{Half-life to stationarity} = \ln(2)/\kappa$$

In our case

$$\ln(2)/0.45 \approx 1.54 \text{ yrs} \cdot \text{the OU half-life}$$

Numerical solution: finite difference

How to solve FP equation on a computer.

Discretize

L-grid: 1000 nodes from L_{\min} to L_{\max}

Time stepping

Crank-Nicolson (implicit · stable)

Sparse linear algebra

Tridiagonal system at each time step

Boundary handling

Reflecting BC at L_{\min} , L_{\max} for OU stability

Verification

Check $\int p \, dL = 1$ throughout

Long-time check

$p(L,t) \rightarrow N(\bar{L}, \sigma^2/(2k))$ as $t \rightarrow \infty$

FP and HJB: the duality

Forward and backward equations are mathematical duals.

Backward HJB

$$\partial V / \partial t + \sup_{\alpha} [b \cdot \partial V / \partial L + \frac{1}{2} \sigma^2 \cdot \partial^2 V / \partial L^2 + r] = 0$$

Forward FP

$$\partial p / \partial t + \partial / \partial L [\alpha^* \cdot p] - \frac{1}{2} \sigma^2 \cdot \partial^2 p / \partial L^2 = 0$$

Duality

FP operator is adjoint of HJB's generator

Why duality matters

Solve HJB \rightarrow know α^* ; plug into FP \rightarrow know μ ; check consistency

Computational use

Forward-backward iteration • standard MFG algorithm

Master equation

Single equation combining both

Extensions: non-OU L processes

What if L doesn't follow OU?

Jump-diffusion L

FP gets non-local term · integro-PDE

Time-varying coefficients

$\kappa(t)$, $\sigma(t)$ · same FP structure, more numerics

CIR process

Cox-Ingersoll-Ross · always positive · FP still applies

Lévy-driven

Pure jump processes · FP becomes integral equation

Heston-like (two state)

Two coupled FP equations · 2D PDE

For GE-LAV

OU baseline + jump extension for crisis (open research)

Applications beyond GE-LAV

Where FP equations show up elsewhere.

Diffusion in physics

Heat equation = FP with no drift

Population dynamics

Density of biological populations

Plasma physics

Vlasov-FP equation

Asset pricing

Density of stochastic vol; underlying for Heston model

Optimal control

Backward equation for V , FP for state evolution

Cross-disciplinary

Same math, many applications

Bridge to Session 23

FP closes the math of how distributions evolve. Now: convexity.

S22 done

FP equation in our toolkit

Next math bridge

Jensen's inequality and convexity

Why convexity matters

Failure 2 of DCF: ignoring convex discount factors

Session 23 covers

Jensen inequality · why averages mislead

Practical impact

0.8-3.6% bias per year by asset class

Reading

Boyd-Vandenberghe Ch. 3 · convexity refresher